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# Dielectric Characterization of B<sub>n</sub> Phases

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Dielectric data of six different  $B_n$  phases formed by banana-shaped molecules are analyzed and discussed with respect to the known results of X-ray measurements. The reorientation about the long molecular axes is discussed with respect to phase transitions. Parallel and antiparallel correlation of the transversal dipole moments in the short range as well as superstructures are detected. Additional AFM and calorimetric investigations are performed. The results of  $B_3$  and  $B_4$  phases prove the existence of dielectric reorientation. Therefore, both modifications cannot be regarded as "ordinary" solids. The existence of the  $B_8$  phase was confirmed.

Keywords: Banana-shaped molecules; mesophases; dielectric investigations; superstructures; AFM

#### INTRODUCTION

Till now a lot of banana-shaped molecules were synthesized and different mesophases of the  $B_n$  type were detected [1-3]. Most of these are characterized and classified by X-ray measurements [4]. Nevertheless

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there are a many open questions related to the order in the short range and the formation of polar superstructures which can be better solved by dielectric spectroscopy and atomic force microscopy [5,6].

#### **EXPERIMENTAL RESULTS**

#### **Experiments**

Dielectric investigations were carried out in the frequency range from 1 Hz to 10 MHz using the Solartron-Schlumberger Impedance Analyzer Si 1260 and a Chelsea Interface. A brass cell coated with gold (d = 0.2 mm) was used as capacitor which was calibrated with cyclohexane. The samples could not be oriented by external magnetic or electric fields. The experimental data of  $\epsilon$ ' and  $\epsilon$ '' were fitted together to the real and imaginary part of equation (1) consisting of two Cole-Cole mechanisms (term 2 and 3), a conductivity contribution (term 4) and term 5 for the description of the capacitance of the double layer at low frequencies

$$\varepsilon^* = \varepsilon_2 + \frac{\varepsilon_0 - \varepsilon_1}{1 + \left(j\omega\tau_1\right)^{1-\alpha_1}} + \frac{\varepsilon_1 - \varepsilon_2}{1 + \left(j\omega\tau_2\right)^{1-\alpha_2}} + \frac{jA}{f} + \frac{B}{f^N}$$
 (1)

with  $\varepsilon_i$  low and high frequency limits of the dielectric constant,  $\omega=2\pi f$  (f - frequency),  $\tau$  - relaxation times ( $\tau=(2\pi f_R)^{-1}$  with  $f_R$  relaxation frequency,  $\alpha$  - Cole-Cole distribution parameters, the conductivity term A ( $\sigma=2\cdot A\cdot \pi\cdot \varepsilon^0$  (specific conductivity),  $\varepsilon^0=8.85\cdot 10^{-12}$  As/Vm) as well as B and N as further fit parameters responsible for the capacity of the

double layer at the interface sample/gold. For the given summary data of different samples are discussed with respect to known X-ray results.

#### The B<sub>1</sub> phase

Fig. 1 shows dielectric constants of the banana-shaped sample

Cr 392 B<sub>1</sub> 439 I (temperatures in K)

in the isotropic (438 K) and the  $B_1$  state (428 K) synthesized at first by Tschierske [7]. There are practically constant  $\varepsilon$ -values by changing the frequency in isotropic phase (if the increase of  $\varepsilon$ ' at low frequencies

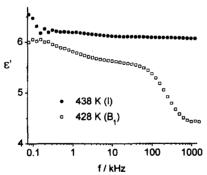


FIGURE 1: Dielectric dispersion of  $\underline{1}$  in the isotropic and the  $B_1$  state.

due to the double layer at the interface gold/sample is neglected) and two dispersion ranges at about 1 kHz and 200 kHz in the B<sub>1</sub> phase. The

high frequency process is related to the reorientation of the molecules about the long axes, the low frequency process probably to a collective motion [5,6]. It should be pointed out that the dielectric constants in the  $B_1$  phase do not exceed the  $\varepsilon$ -values in the isotropic phase at all temperatures. This points to an antiparallel correlation of the longitudinal dipole moments of  $\underline{1}$  and it is in agreement with the structure model of a rectangular columnar phase proposed at first by Sekine et al. [4].

#### The B, phase

Dielectric results on the  $B_2$  phase are already reported in the literature [6,8-10]. Similar to the dielectric behavior of the  $B_1$  modification low and high frequency relaxation were detected [8, 9]. Additionally a very fast reorientation of terminal dipoles could be seen at low temperatures [6]. Experimental data on sample  $\underline{2}$ 

are shown in Figure 2. Contrary to the  $B_1$  phase, an increase of the dielectric constant to about 10 at 10 kHz was measured in the  $B_2$  modification.  $\epsilon$ ' decreases strongly at the I/B<sub>2</sub> transition. At frequencies of about 200 kHz the dispersion due to the reorientation about the long axes reduces the  $\epsilon$ '-values. This process could be also detected in the isotropic phase, but with a stepwise increase of the relaxation time and

a much lower intensity [6, 8]. Thus, this absorption at about 2 MHz cannot clearly be seen in Figure 2. We explain the increase of the

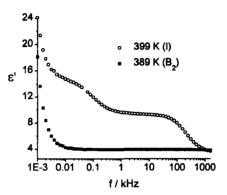


FIGURE 2: Dielectric dispersion of 2 in the isotropic and the B<sub>2</sub> state.

absorption intensity at the  $I/B_2$  transition as result of the ferroelectric order in the short range. The relaxation frequencies  $f_R$  for the reorientation about the long molecular axes (open circles) are also strongly influenced by the phase transition  $I/B_2$  as seen in Figure 3.

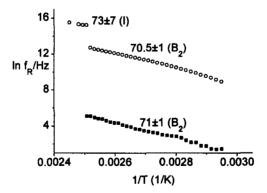


FIGURE 3: Relaxation frequencies of **2** versus the reciprocal temperature. The numbers are activation energies in kJmol<sup>-1</sup>.

Measurements on sample  $\underline{2a}$  (similar to  $\underline{2}$ , only the terminal  $-C_{14}H_{29}$  are substituted by polar octyloxy groups) result in much higher limits for the high frequency dielectric constants  $\varepsilon_2$  (see Figure 4). The octyloxy groups show a separate relaxation in the GHz range. This means that the ferroelectric short range order in  $B_2$  is produced only by the stiff middle parts of the molecules.

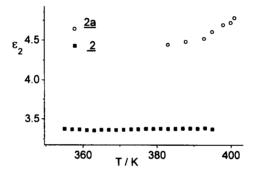


FIGURE 4: High frequency limit of ε' of 2 and 2a.

The low-frequency absorption, not seen in the isotropic phase, is connected with a collective dynamics within the antiferroelectric layers. In order to prove the existence of a superstructure AFM measurements were carried out. An AFM image obtained on a sample which could be supercooled [11] is presented in Figure 5. At lower magnifications a typical focal-conic texture with additional disclinations between the minima at the surface is observed. In line measurements of Figure 5 steps in the surface profile of one layer are seen. In other samples additionally steps of two and three layers were detected. From these results

it can be concluded that there are not very strong interactions between two neighboring layers with antiferroelectric dipole order. Helical super-structures could not be observed.

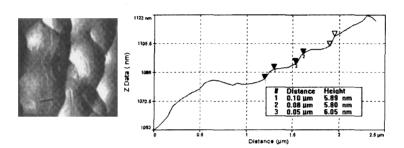


FIGURE 5: AFM-image of a sample in the glassy B<sub>2</sub> state. The line inside of the image represents the part for the line measurements shown right. Steps in the profile are indicated and given in the inset.

### The B<sub>3</sub> phase

Sample 3 is a classical "banana-shaped liquid crystal"

$$C_{14}H_{29}$$

I 433.4 B<sub>2</sub> 396 B<sub>3</sub> 330 Cr (by cooling)

synthesized at first by Sekine and coworkers [12]. The B<sub>3</sub> phase exhibit many sharp X-ray pattern. Therefore, it was classified as solid modification [4]. Dielectric constants of all phases are presented in Figure 6. Let us take the dielectric constants in the isotropic phase (solid squares)

as reference for the discussion. Again much higher  $\varepsilon$ '-values in the B<sub>2</sub> phase with the MHz dispersion are detected (open triangles). The low-frequency relaxation at 300 Hz is superimposed by the double layer due to the high conductivity.

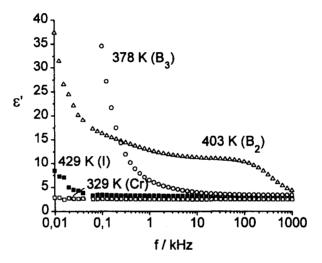


FIGURE 6: Dielectric constants of 3 in different phases.

The solid phase (open squares) shows the lowest dielectric constant. Here all motions are frozen. In the  $B_3$  phase (open circles) higher dielectric constants than in the isotropic state were found. Furthermore a strong increase of  $\varepsilon$ ' at low frequencies is detected which results from a strong increase of conductivity [13] and the corresponding double layer at the interface capacitor/sample. The molecular origin for the increase of conductivity is difficult to understand. The small difference in  $\varepsilon$ ' between the solid, the isotropic and the  $B_3$  phases points to a small

net dipole moment of sample <u>3</u>. The higher ε'-value in the B<sub>3</sub> phase may be the result of a local dynamics like in plastic crystals. Thus, one can understand the appearance of many sharp X-ray reflections [4]. An AFM-image of sample <u>3</u> is presented in Figure 7. To obtain this image a thin layer of the sample was heated to 430 K, cooled to 422 K with a rate of 0.5 Kmin<sup>-1</sup> then cooled with 5 Kmin<sup>-1</sup> to 400 K and afterwards very fast to room temperature. Surprisingly, the AFM-picture

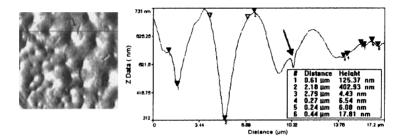


FIGURE 7: AFM-image of a sample in the B<sub>3</sub> state. Along the line inside of the image the line measurement shown right was done. Breaks are indicated by arrows. Steps in the profile are related to the minima of the focal conics (#1-3) and to distances of one layer (#4-6).

remembers on that of the B<sub>2</sub> phase given in Figure 5. There is a focal conic texture, but the number of disclination lines between the minima is reduced. Furthermore there are cracks in the surface profile. These points to a layered structure deformed due to volume contraction at the phase transition. Probably we are dealing with an epitaxy effect from the B<sub>2</sub> phase, but in any way regular crystals cannot be detected.

#### The B<sub>4</sub> phase

Dielectric investigations on

I 444.3 B<sub>2</sub> 425.2 B<sub>3</sub> 412 B<sub>4</sub> (by cooling)

are presented in Figure 8. The dielectric constants of the isotropic (solid squares),  $B_2$  and  $B_3$  phases show similar behavior as already presented in Figure 6. The solid state could not be reached, thus there is no chance for a direct comparison as done in Figure 6. Nevertheless

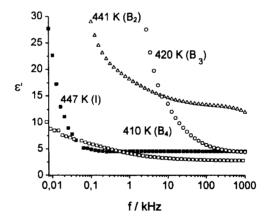


FIGURE 8: Dielectric constants of 4 in different phases.

it is obvious, that the  $\epsilon$ '-data in the  $B_4$  phase (open squares) cross those of the isotropic state. The strong increase of the data in the isotropic

phase below 50 Hz and in  $B_4$  below 10 Hz is connected with the double layer. This effect could be separated from a low-frequency absorption in  $B_4$  ( $f_R$  about 200 Hz). X-ray pattern of the  $B_4$  phase show wide angle reflections with a broad distribution pointing to fluctuations in the lateral distances [4]. Thus, the observed low-frequency absorption may be related to a complete or partial reorientation about the long molecular axes of 4.

#### The B<sub>2</sub> and B<sub>8</sub> phases

The  $B_5$  and  $B_6$  phases are not well characterized till now. Interesting results were obtained on the  $B_7$  phase of

Cr 393 B<sub>7</sub> 457 I

In this case the dipole moments compensate each other in the middle part. Therefore the reorientation about the short axes is difficult to detect. On the other hand the experimental data of 5 can be well fitted under assumption of a low-frequency absorption as it was demonstrated in [5]. The calculated static dielectric constants show a maximum at about 400 K. The directly measured dielectric constants outside of the relaxation range (60 kHz) at cooling and heating are presented in Figure 9. Astonishingly the  $\epsilon_0$ -values show a step at about 400 K and

seem to be stable up to the beginning of crystallization. This effect could be reproduced by heating. Calorimetric and AFM-measurements support an interpretation as phase transition.

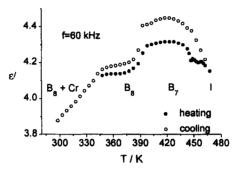


FIGURE 9: Measured dielectric constants of 5.

#### **SUMMARY**

At the present stage of investigations one can summarize the results in the following way:

- 1. In the B<sub>1</sub>, B<sub>2</sub>, and probably also in the B<sub>4</sub> and B<sub>7</sub> phases the molecules can reorient about their long axes.
- In the B<sub>1</sub> modification an antiparallel dipolar correlation and in B<sub>2</sub> an ferroelectric short range order was detected.
- In all phases a restriction of the dynamics with respect to that of the isotropic state is seen.
- 4. In the B<sub>3</sub> and B<sub>8</sub> phases a clear restriction of the dynamics and a reduction of the dielectric increments was observed. It is possible that only parts of the molecules can rotate in the investigated time

- scale. On the other hand the B<sub>3</sub> phase shows an anomalous high conductivity and focal conic domains at the surface.
- 5. Low-frequency relaxation which are related to super-structures were detected in the B<sub>1</sub>, B<sub>2</sub>, B<sub>4</sub>, and B<sub>7</sub>, phases.

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